



ELECTRON-MICROSCOPIC INVESTIGATION OF A MODULATED (TWEED) STRUCTURE AND OPTICAL CONDUCTIVITY IN ALLOYS AND Y-Ba-Cu-O SINGLE CRYSTALS

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A correlation between the structure-unstable state of nonstoichiometric $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ oxides and the loss of metallicity has been established. A comparison with similar states in metal systems is carried out.

Some metallic alloys, such as Ti-Nb, Ti-V, Ni-Al, are known to possess a correlation between optical properties untypical for metals and a specific structure state of these alloys. This state arises before a phase transition either in a critical concentration range, or near it. Thus, for Ti-V alloys the concentrations of 13 - 15 at. % V are the critical ones, and on quenching from 1000 °C the BCC β -phase transforms into ω -phase with atomic displacements from the points of a hexagonal lattice. In that state electron diffraction patterns exhibit not only the reflections of ω -phase, but also an intensive diffuse scattering along the $\{111\}^*$ octahedral planes resulting from atomic displacements.

At some deviation from the critical concentration in Ti-20 at. % Nb alloy an even more intensive diffuse scattering is observed. It has a form of hollow figures inscribed

into octahedrons with planes along $\{111\}^*$, and with intensity maximums in the positions near the nodes of the ω -phase lattice (an incommensurable ω -phase, Fig. 1.). In this case a tweed contrast corresponding to the atom dis-

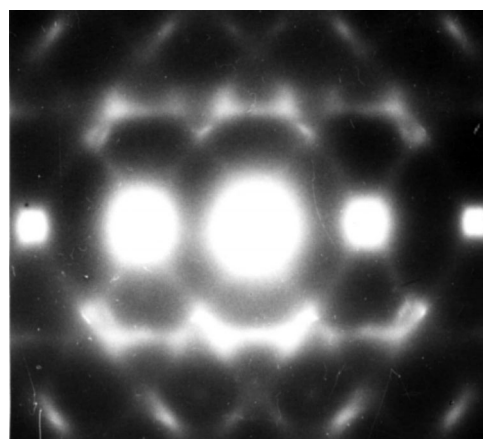


Fig. 1. Electron diffraction pattern of alloy Ti-20 at.% V

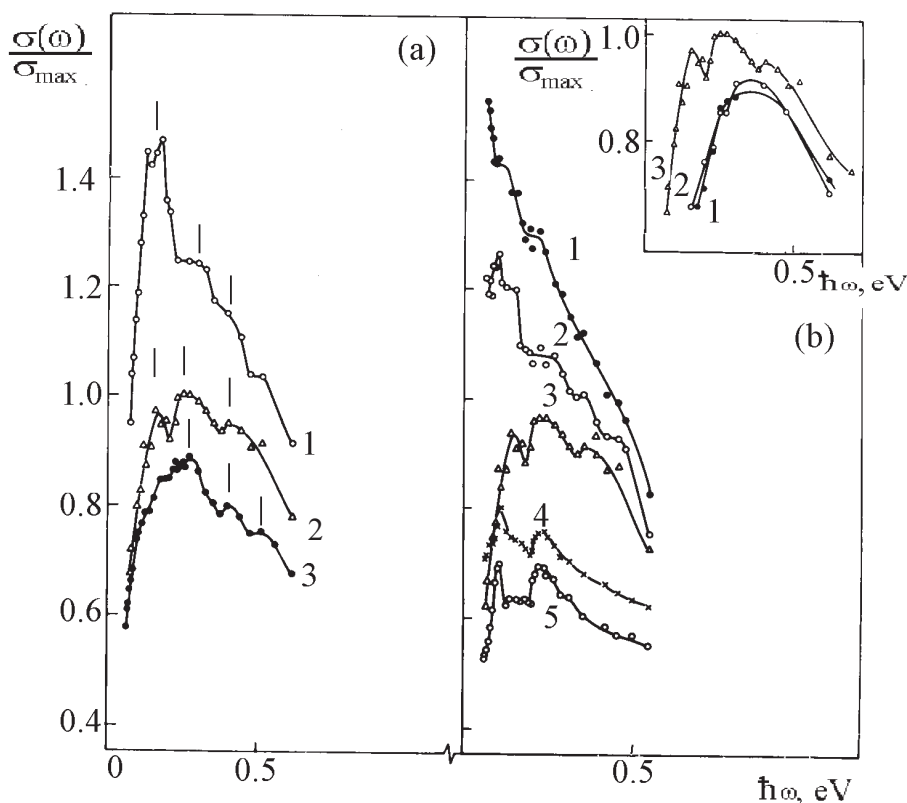


Fig. 2. Optical conductivity of alloys :

a) - Ti-25 at.% Nb (1), Ti-20 at.% V (2), Ti-8 at.% V + 4 at.% Mo (3); b) - Ti-50 at.% V (1), Ti-30 at.% V (2), Ti-20 at.% V (3), Ti-15 at.% V (4), Ti-13 at.% V (5). Inset shows interband contribution to $\sigma(\omega)$ of Ti alloys with 30 at.% V (1) and 50 at.% V (2) and experimental $\sigma(\omega)$ of alloy Ti-20 at.% V (3)

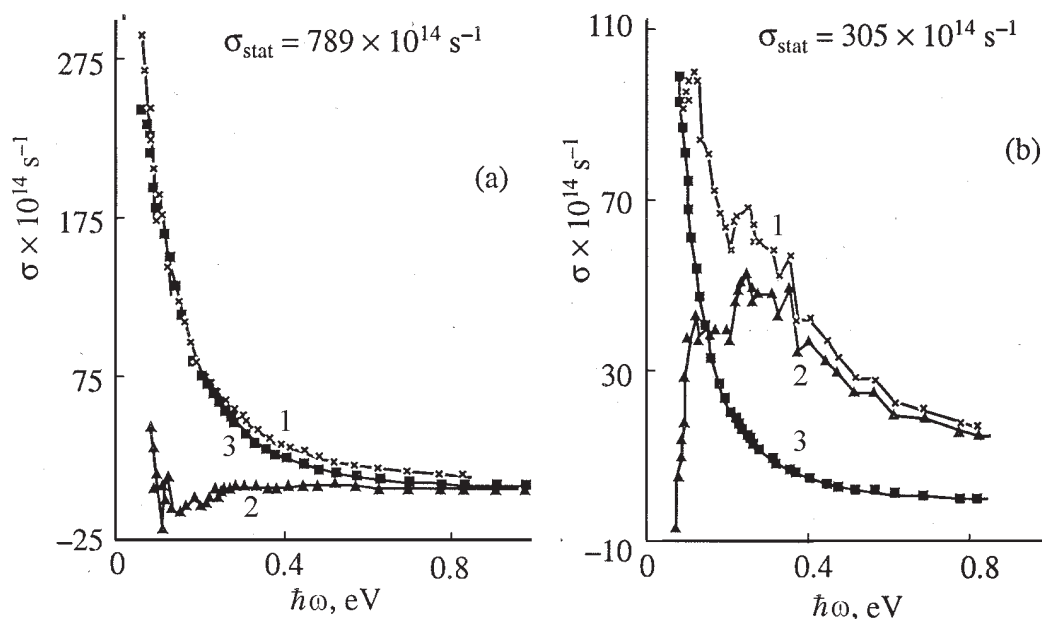


Fig. 3. The dispersion of optical conductivity $\sigma(\omega)$ in Ni-Al alloys with 50 at.% Ni(a), 55 at.% Ni(b): 1-experiment; 2-contribution from interband transitions; 3-calculation by Drude-Ziner equation

placement waves and the space ordering of ω -phase nuclei is observed on micrographs of the structure.

On the $\sigma(\omega)$ plots of Ti alloys with 13, 15, 20 at.% V there are two pronounced singularities: an absorption band with a maximum at 0.25 eV and the absence of a rise from intraband absorption (i.e. the absence of the Drude-Ziner contribution from free carriers) down to energy 0.07 eV (Fig. 2). Nevertheless, the relation $k(\lambda) > n(\lambda)$ indicating the metallic nature of conductivity, is satisfied over the whole investigated infrared (IR) range of the spectrum. Optical conductivity $\sigma(\omega)$ of the Ti alloys with 30 and 50 at.% V has the usual frequency dependence for metals in the IR range [1]. With an increase of vanadium content to 30 at.% the diffuse scattering intensity markedly decreases. The transition from the anomalous optical properties to the normal ones is observed at vanadium content above 20 at.%. It is interesting to note that the alloys can be divided into the same two groups (with anomalous and normal optical properties) according to the sign of the temperature coefficient of resistivity [2], whilst the alloys with a negative coefficient exhibit unusual dispersion of the optical parameters.

An abrupt drop of the square of the plasma frequency is observed on its concentration dependence in the Ti-V alloys with 13-20 at.% V [1]. Thus, the unusual optical properties are found in heterogeneous alloys with unstable structure and ω -regions possessing large atomic displacements and forming a macrolattice in the crystal.

High-frequency conductivity dispersion curves for Ni-Al alloys, as measured in the range 0.07-1.22 eV, are presented in Fig. 3. While the conductivity curve for the alloy with 50 at.% Ni is typical for frequency dependence of metals in the IR region, the $\sigma(\omega)$ of alloys with 55 and 59 at.% Ni has an anomalous behaviour of their optical conductivity: with energy decrease $\sigma(\omega)$ at first increases monotonously, and then such a Drude-like growth is disturbed by a slight fall in a low energy region starting from the en-

ergy $h\nu = 0.12$ eV. As a result, an absorption band is revealed around 0.2 eV [3].

The alloys with 55 and 59 at.% Ni (B2 matrix) are characterized by a planar diffuse scattering on $\{111\}^*$ planes and rather strong rods along $\langle 110 \rangle^*$ directions which extinguish according to the law for transverse polarized waves and correspond to shear displacements of the $\{110\}$ planes along the $\langle 1\bar{1}0 \rangle$ directions. Electron micrographs reveal a tweed contrast in these alloys, especially strong in the 62.5 at.% Ni-alloy. The electron diffraction patterns of the 62.5 at.% Ni alloy demonstrate the most pronounced effects of the diffuse scattering (Fig. 4).

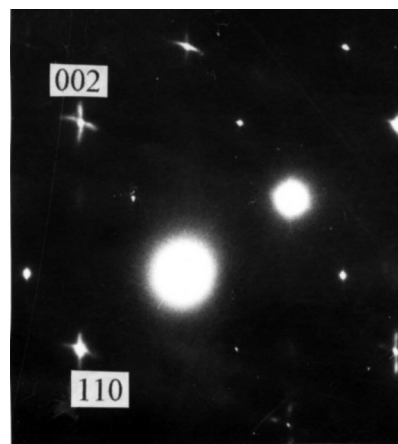


Fig. 4. Electron diffraction pattern of Ni-Al alloy with 62,5 at.% Ni

Now the similar structure and the metallicity loss have been found in the structure-unstable oxides $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$. According to the Khachatryan and Morris theory [4], the decomposition occurs only in the nonstoichiometric compounds with $\delta > 0.2$, while the oxygen-saturated samples do not decompose. We must verify this experimentally.

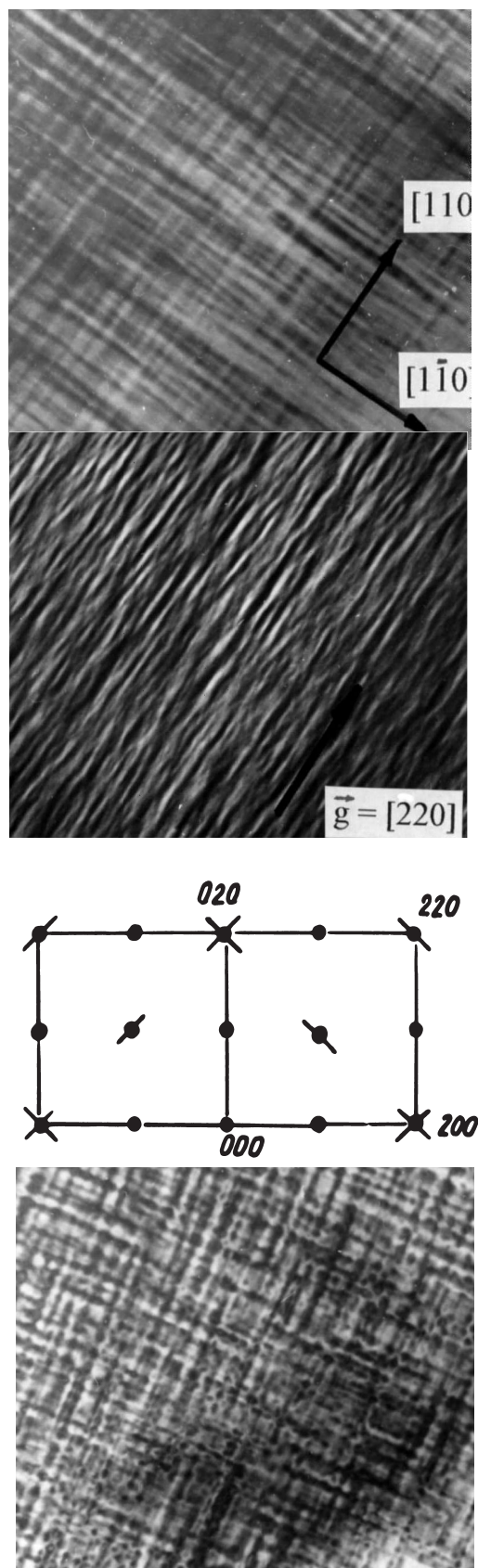


Fig. 5. Structure of a nonstoichiometric $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ single crystal. Magnification $\times 146000$.

- (a) bright-field image; $\delta \approx 0.2$; after annealing at 200°C for 100 h;
 (b) dark-field image using 220 reflection;
 (c) scheme of an electron diffraction pattern, zone axis $[001]$; $\delta \approx 0.2$; after annealing at 200°C for 100 h;
 (d) bright field image; $\delta \approx 0.6$; after annealing at 910°C for 24 h.

Therefore, in the present work we paid special attention to the influence of low-temperature prolonged annealing (at 200°C for 100 h) on the fine structure and optical properties of nonstoichiometric single crystals.

By electron-microscopy investigation and magnetic susceptibility measurements it has been found that in progressive concentration transition from $\text{YBa}_2\text{Cu}_3\text{O}_7$ to $\text{YBa}_2\text{Cu}_3\text{O}_6$ phase a wide region of structure instability appears, accompanied by the occurrence of periodic and aperiodic waves. The 123 system in the intermediate region of oxygen concentrations is a superposition of two orthorhombic phases with different parameters and oxygen content. The instability of $\text{YBa}_2\text{Cu}_3\text{O}_{6.8}$ compound is found only during a low temperature annealing at 200°C for 100 h. An early stage of the decomposition with the formation of tweed structure is observed in that case. Besides fine (3-5 nm) precipitates of an oxygen-enriched phase, transversely polarized waves of atomic displacements in the $[1\bar{1}0]$ and $[110]$ directions arise in the oxygen-depleted matrix, producing the observed tweed contrast (Fig. 5a). When the diffraction conditions change and the dark-field image is produced in (220) reflection, one of the systems of bands disappears, and the other one is parallel to the \mathbf{g}_{220} vector (Fig. 5b). In the electron-diffraction patterns the diffuse scattering effects are seen as rods along the $[110]^*$ directions crossing the nodes of a reciprocal lattice, with extinction typical for transversely polarized waves (Fig. 5c).

In the $\text{YBa}_2\text{Cu}_3\text{O}_{6.4}$ single crystals decomposition into two phases takes place already during cooling from 900°C because of the smaller stability of the compound with lower oxygen content. A pronounced tweed structure is observed consisting of a grid of orthorhombic phase precipitates and the two systems of transversely polarized atomic displacement waves in the $[110]$ and $[1\bar{1}0]$ directions (Fig. 5d).

Optical properties of $\text{YBa}_2\text{Cu}_3\text{O}_{6.8}$ single crystals in an initial state and after a low-temperature annealing at 200°C , 100 h, were investigated in the spectrum range of 0.5 - 5 eV at 290 K. The optical constants $n(\lambda)$ and $k(\lambda)$ of single crystals of the 123 system were measured by the ellipsometric technique. The relative experimental errors for $n(\lambda)$ and $k(\lambda)$ were 2 % - 4 %. The automated ellipsometer for measurements on small samples was assembled on the basis of the KSVU-12 spectrometer.

Optical conductivity spectra $\sigma(\omega)$ of $\text{YBa}_2\text{Cu}_3\text{O}_y$ compounds are already known and described elsewhere [5]. Oxygen depleted specimens with $y = 6.1$ are known to have three peculiarities: an absorption peak at 1.7 eV associated with a fundamental edge; a wide band at 2-3 eV and an intensive peak at 4.1 eV. The $\sigma(\omega)$ spectrum of specimens with $y = 6.9$ is modified in the following way. A wide band at 2 - 3 eV is retained, the peak at 4.1 eV disappears, and the peak at 1.7 eV is veiled by a contribution of free carriers dominating in the spectrum range of $E < 1.0$ eV. Fig. 6 presents $\sigma(\omega)$ curves of $\text{YBa}_2\text{Cu}_3\text{O}_{6.8}$ specimens before and after the annealing. It should be mentioned that in a spectrum of an initial specimen a weak singularity near 4.0 eV inherent to oxygen depleted specimens has been found. That has been also testified by an absence of a dramatic $\sigma(\omega)$ growth for energies $E < 1$ eV. Radical changes in the

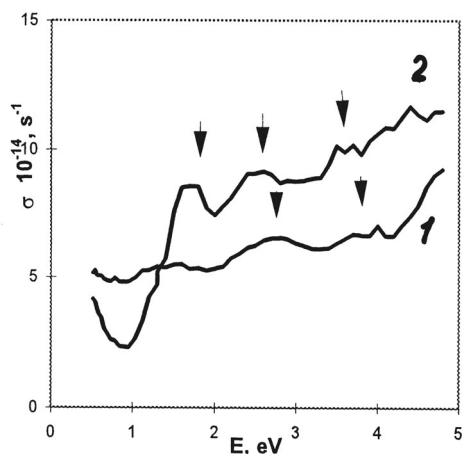


Fig. 6. Optical conductivity $\sigma(\omega)$ of $\text{YBa}_2\text{Cu}_3\text{O}_{6.8}$ before annealing (1) and after annealing (2).

$\sigma(\omega)$ spectra are observed after the annealing. In general the behaviour of $\sigma(\omega)$ differs from that of not annealed specimens both with $\gamma = 6.1$ and with $\gamma = 6.9$. The annealed specimen spectra demonstrate features of both oxygen depleted state and metal state, namely, the sharpening of an absorption edge and a peak at 1.7 eV, along with a presence of an appreciable contribution of free carriers have been observed. Nevertheless, a decrease of free carrier contribution in the annealed state in comparison with the initial one is clearly seen when we analyse a spectrum of a characteristic electron energy losses function. After the annealing a screened plasma frequency $\bar{\omega}_p = \omega_p/\epsilon_\infty$ (where ω_p is a plasma frequency of current carriers and ϵ_∞ is a bare dielectric permeability) decreases from 1.65 eV to 1.1 eV.

Thus, the matrix oxygen depletion of the annealed specimen is manifested in a substantial decrease of metallicity and an appearance of a peak at 1.7 eV in the absorption spectrum, typical of specimens with a high oxygen deficiency.

Thus, for a number of different materials a correlation between the structure-unstable state and metallicity loss has been found.

References

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