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## X-RAY DIFFRACTION SYSTEMS – TECHNICAL ADVANTAGES AND APPLICATION RELATED BENEFITS

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Experimental factors such as crystal characteristics, available experiment time and the properties of the X-ray sources and detectors have a strong impact on data quality and can make the difference between success and failure in phasing attempts or result in a more or less accurate atomic model. The talk focuses on the XRD equipment and is intended as an overview on recent developments of X-ray diffraction systems and the fields where STOE XRD instruments are most beneficial.

A variety of measurement setups with respect to goniometers, diffraction geometries, detectors, X-ray

sources and sample environments will be presented, and both Powder XRD and Single Crystal XRD applications will be exemplified. Additionally, recent advances in detector technology, novel in situ camera and implementation of MetalJet X-ray source will be highlighted and the tangible benefits for the scientists will be made transparent, e.g. gaining measurement speed, improving data quality and acceptance of samples with complex crystallinity.

### Session V, Tuesday, June 19

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## STRUCTURAL DESCRIPTION AND PROPERTIES OF Mg<sub>2</sub>Al-LAYERED DOUBLE HYDROXIDES INTERCALATED WITH THE FLUVASTATIN ANIONS SOLVED BY MOLECULAR SIMULATION METHODS

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## IN-SITU SANS STUDY OF PRECIPITATES NANOSTRUCTURE OF SINGLE CRYSTAL Ti-15Mo

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Titanium alloys have plenty of applications in industry and medicine due to unique combination of high strength, low density, and excellent biocompatibility [1]. Ti-15Mo (in wt.%) is called  $\beta$ -stabilized binary alloy, it contains mostly  $\beta$ -phase (bcc) and also encloses metastable precipitates of  $\alpha$  (hexagonal) and  $\alpha'$  (hcp) phases [2]. Microstructure of the precipitates has great impact on mechanical properties and thermal stability of the alloy. Thanks to nanometres size of the precipitates and partitioning of molybdenum content in different phases small-angle neutron scattering (SANS) instrument can be effectively used for the in-situ investigation of this microstructure.

In the present studies SANS data were measured at three orientations of the single crystal sample – (111), (110) and (100) of  $\beta$ -phase with correspondent plane perpendicular to incident neutron beam direction. Samples

were installed in vacuum high temperature furnace and heated with heating rates of 1 K/min from room temperature up to 600 °C. SANS data were recorded in so-called list-mode and afterwards binned by time frames of 5 minutes, which corresponds to temperature range of 5 K. The measured data were calibrated using water and corrected by standard measurements of cadmium background. Scattering of the sample in high temperature furnace was used as “buffer” background.

2D pattern of SANS for sample orientation [100] parallel to incident neutron beam at temperature range 410 °C ÷ 415 °C taken at sample-to-detector distance SD = 12 m with collimated neutron beam of 5 Å (± 0.5 Å) wavelength is shown in Figure 1. The observed reflexions are formed by interparticles structure factor, due to high ordering of  $\alpha$ -precipitates. This ordered microstructure exists in wide



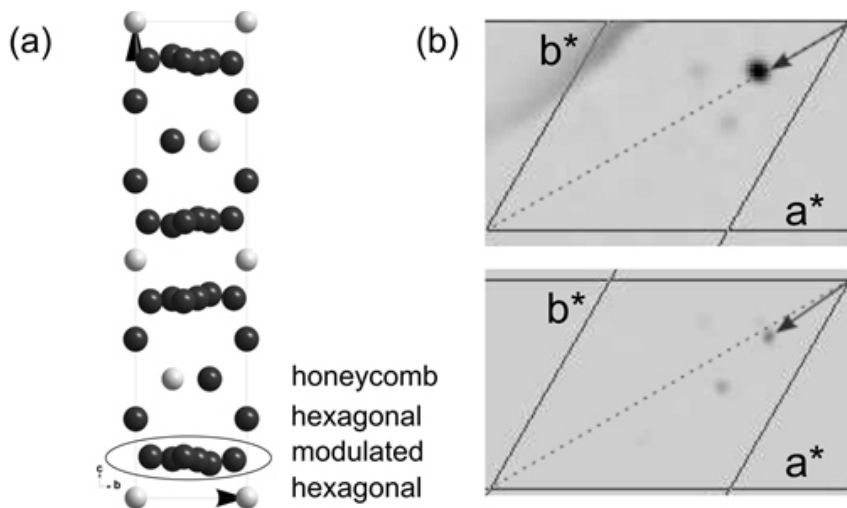


studied and for  $\beta$  and  $\gamma$ . Although the point group of the complete diffraction pattern is  $\bar{3}$ , the average structure can be described in the space group  $P6_3/mmc$ . The average structure can be described by three repetitions of hexagonal and honeycomb shaped layers (Fig. 1a). The modulation is revealed by large anisotropic distribution of electron density around the Cu atoms of the layer B. Two modulation vectors are necessary to index the complete diffraction pattern (Fig. 1b):  $q_1 = (0, 0, 1/3)$  and  $q_2 = (0, 0, 1/3)$ , where  $\delta = 0.2509(10)$  (point group  $\bar{3}m$ ) for  $\beta$ , and  $\delta = 0.23458(7)$ ,  $\epsilon = 0.28171(7)$  (point group  $\bar{3}$ ) for  $\gamma$ . Both structures could be solved in (3+2)D superspace by Superflip [4], superspace groups  $P\bar{3}1c(0, 0, 1/3)(0, 0, 1/3)$  and  $P\bar{3}(0, 0, 1/3)(0, 0, 1/3)$  for  $\beta$  and  $\gamma$ , respectively. The modulation function of both structures have amplitudes comparable to the size of the unit cell, discontinuities and windows. Given the complexity of the modulation and the small number of reflections measured, the function could not be parametrized to be refined using the superspace formalism, and a supercell approximation had to be used. Since  $\delta \sim 1/4$  for  $\beta$ , a  $4 \times 4 \times 3$  supercell could be used, while the smallest possible approximation for  $\gamma$  was a  $14 \times 14 \times 3$  supercell. The refinement was performed in Jana2006 [5].

Temperature-dependent powder X-ray diffraction of the samples  $Cu_{74}Si_{26}$  and  $Cu_{78}Si_{22}$  was measured every  $30^\circ$ , from  $30^\circ C$  to  $700^\circ C$  with heating rate of  $5^\circ C/min$ . Two cycles of heating and cooling were measured to verify the reversibility of the transitions. Three additional phases, which were not present in the phase diagram, were observed, the transitions were reversible and reproducible

with small hysteresis in the transition temperatures. Le Bail fitting of the powder patterns was performed in Jana2006 using the models obtained by SCXRD, pseudo-Voigt profile, manual background combined with fifteen terms of Legendre polynomials. Initially the main reflections were indexed using cyclic refinement, and the modulation vectors were refined separately for each temperature, after the cyclic refinement. Except for the phase  $\beta$ , powder patterns of the sample  $Cu_{78}Si_{22}$  presented the same transitions as for the sample  $Cu_{74}Si_{26}$  and only this sample will be shown. From the six phases observed in the sample  $Cu_{74}Si_{26}$ , five were completely indexed. Our study shows that the phase diagram might be more complex than that reported in the literature.

1. Olesinski, R.W., Abbaschian, G.J., Bull. Alloy Phase Diagrams 7, 170 (1986).
2. Corre a C.A. et al., Acta Crystallographica B73, 767-774 (2017).
3. Corre a C.A. et al., Intermetallics 91, 129-139 (2017).
4. Palatinus, L. and Chapuis, G., J. Appl. Cryst. 40, 786-790 (2007).
5. Petri ek V., Du ek M., Palatinus L., Z. Krist. 229 (5) 345-352 (2014).



**Figure 1.** (a)  $Cu_3Si$  structure viewed along  $a$  - Cu - black, Si - grey. Strongly modulated honeycomb layer filled with copper causes two variants of the structure shown in (b) above, modulation vectors for  $\beta$   $q_1 = (0, 0, 1/3)$  and  $q_2 = (0, 0, 1/3)$ , where  $\delta = 0.2509(10)$ , and below, for  $\gamma$   $\delta = 0.23458(7)$ ,  $\epsilon = 0.28171(7)$ .

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## TRANSITIONS TOWARD COMPLEX ELECTRONIC STATES AND SUPERPERIODIC STRUCTURES IN $P_4W_{16}O_{56}$

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The MonoPhosphate Tungsten Bronzes (MPTB) family,  $(PO_2)_4(WO_3)_{2m}$ , can be described by a regular intergrowth of  $(PO_4)$  tetrahedra layers and of slabs constituted by corner-sharing- $(WO_6)$  octahedra, with a thickness depending on the  $m$  parameter. These low-dimensional oxides are known to exhibit successive transitions toward Charge Density Wave (CDW) states. These transitions are associated to lattice distortions leading to the appearance of incommensurate or commensurate structural modulations [1]. In this family, the electronic anisotropy and the density of carriers of the system can be tuned by modifying the thickness of the  $WO_3$  slabs *i.e.* changing  $m$ . MPTB family is thus a relevant system to analyse the effect of the dimensionality on the CDW electronic instabilities. Temperature-dependant X-ray diffraction (XRD) [1] and transport measurements reported in the literature, for different terms of the family, reveal a significant change of behaviour between the terms with a low and high value of  $m$ ,  $m < 7$  and  $m > 7$  respectively. Classical CDW transitions are reported for the low terms, characterized by a smooth resistivity jump and by the formation of clusters of tungsten in the centre of the  $WO_3$  slabs [2]. For the high terms [3], a structural transition is observed in XRD but the electronic transport studies do not show the usual signature attributed to a CDW. Moreover, the only structural study performed on a high term in the modulated state ( $m = 10$ ) [3] evidences

anti-ferroelectric-type (AFE) atomic displacements for the tungsten atoms without reporting of the formation of clusters of tungsten.

We will present both the transport properties and the analysis of the structural modulations for the  $m = 8$  term. Three first-order transitions associated with large thermic hysteresis were identified. The analysis of the structural modulations characterizing the different states, *via* the use of the super-space formalism, reveals the existence of AFE-type displacements and the formation of clusters for the tungsten atoms. These signatures can be assigned to the coexistence of AFE and CDW properties in the material. These two properties are *a priori* incompatible, but an extensive study of the transport properties versus temperature supports this hypothesis. This result enlightens the very interesting position of  $P_4W_{16}O_{56}$  ( $m = 8$ ) in the border area between the low and the high  $m$  values in the MPTB family to discuss the competition regime between CDW and ferroelectric instabilities.

1. A. Ottolenghi, J. P. Pouget, *J. Phys. I France*, **1996**, Vol 6, p. 1059-1083.
2. J. Lüdecke, A. Jobst, S. Van Smaalen; *Europhys.Lett.*, **2000**, Vol.49, n°3, p.357–361.
3. P. Roussel et al., *Phys. Rev. B*, **2000**, Vol.62, p.176.